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Kodak

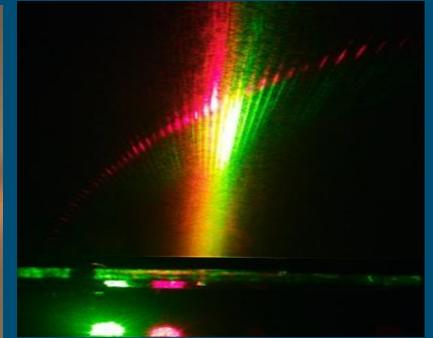
NOTTINGHAM
TRENT UNIVERSITY

NTU

Immersed Superhydrophobicity, Reconfigurable Substrates and Shaped Liquid Surfaces

Glen McHale

School of Science & Technology



www.naturesraincoats.org

22nd February 2011

Overview

1. Concepts of Water Repellency
2. Effects of Topography
3. Amplification, Switching and Superspreading
4. Conformable Substrates
5. Immersed Superhydrophobic Surfaces
6. Dielectrowetting

1. Concepts of Water Repellency

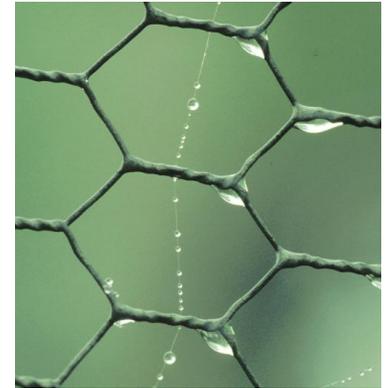
Surface Tension

Liquid Surface

Molecules at a surface have fewer neighbours

Liquid surface ("skin") behaves as if it is in a state of tension

For a free "blob", the smallest area is obtained with a sphere



<http://www.brantacan.co.uk>

Size Matters

Surface tension force \propto length

e.g. Force $\sim R\gamma_{LV}$

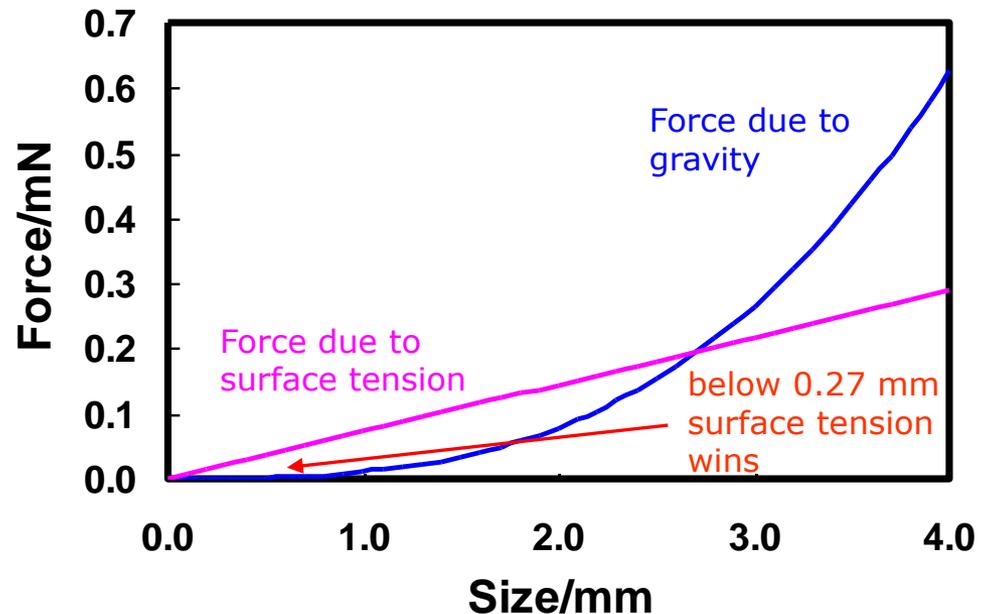
Gravity forces \propto length³

e.g. Force $\sim R^3\rho g$

Small size \Rightarrow surface tension
wins

Small means \ll capillary length

$$\kappa^{-1} = (\gamma_{LV}/\rho g)^{1/2} \sim 2.73 \text{ mm (water)}$$



Size Matters: Fiction or Fact?



The Movie – Antz (1998)

Copyright: DreamWorks Animation (1996)



Courtesy: BigWave Productions



Is it just imagination?
Or could it happen?

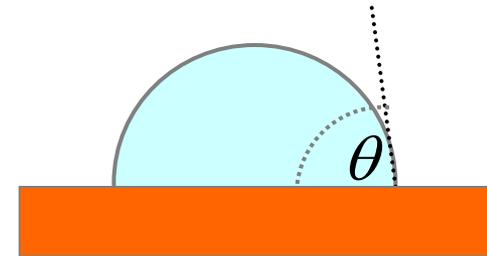
Hydrophobicity and Superhydrophobicity

Surface Chemistry

Terminal group determines whether surface is water-liking or water-fearing
Hydrophobic terminal groups are Fluorine (CF_x) and Methyl (CH_3)

Contact Angles on Teflon

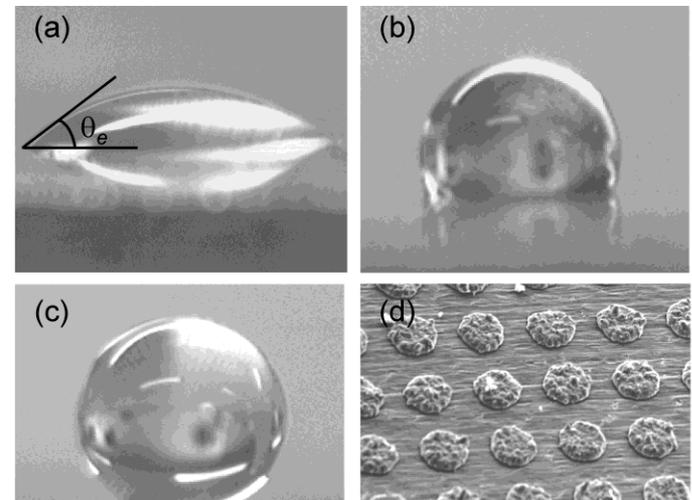
Characterize hydrophobicity
Water-on-Teflon gives $\sim 115^\circ$
The best that *chemistry* can do



Enhancement by Topography

- (a) is water-on-copper
- (b) is water-on-fluorine coated copper
- (c) is a super-hydrophobic surface
- (d) "chocolate-chip-cookie" surface

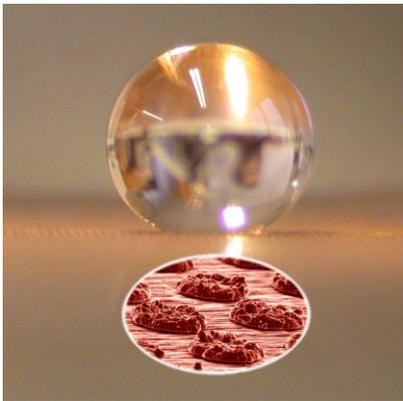
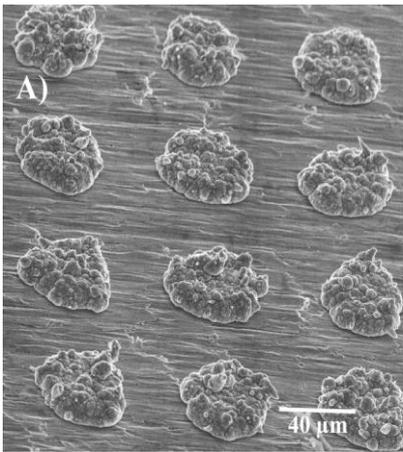
***Superhydrophobicity is when $\theta > 150^\circ$
and a droplet easily rolls off the surface
(low contact angle hysteresis)***



2. Effects of Topography

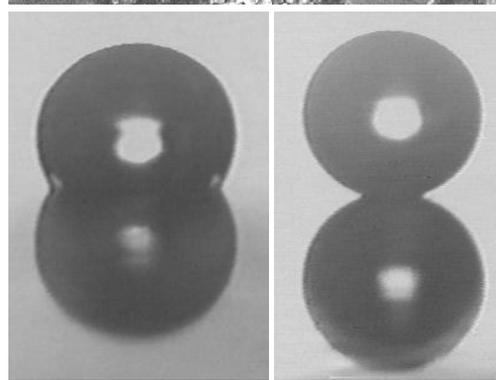
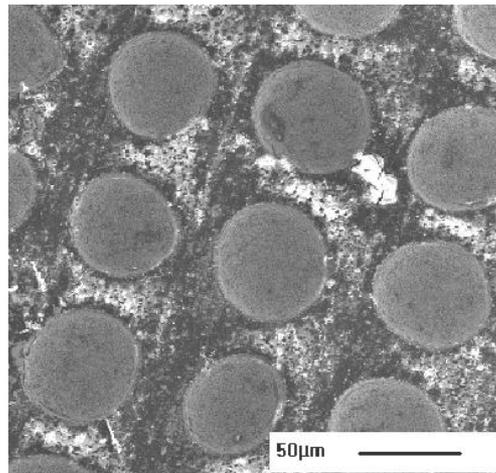
Superhydrophobicity – NTU Examples

Deposited Metal



Patterned & hydrophobic

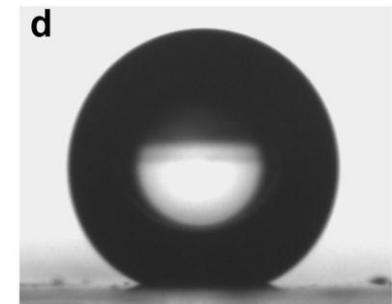
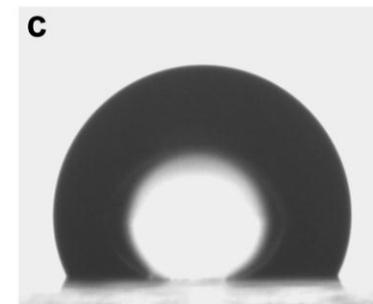
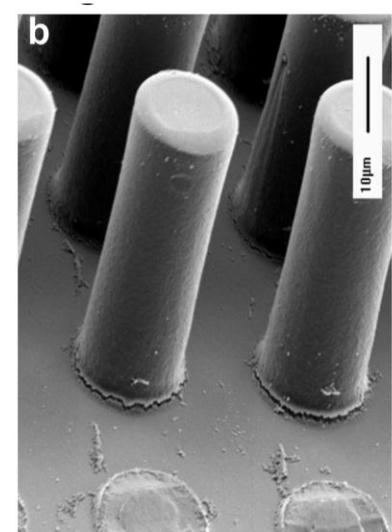
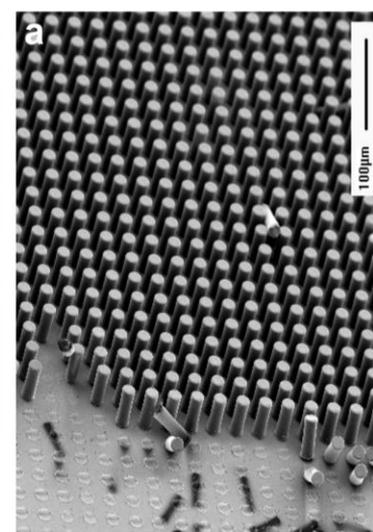
Etched Metal



Flat & hydrophobic

Patterned & hydrophobic

Polymer Microposts

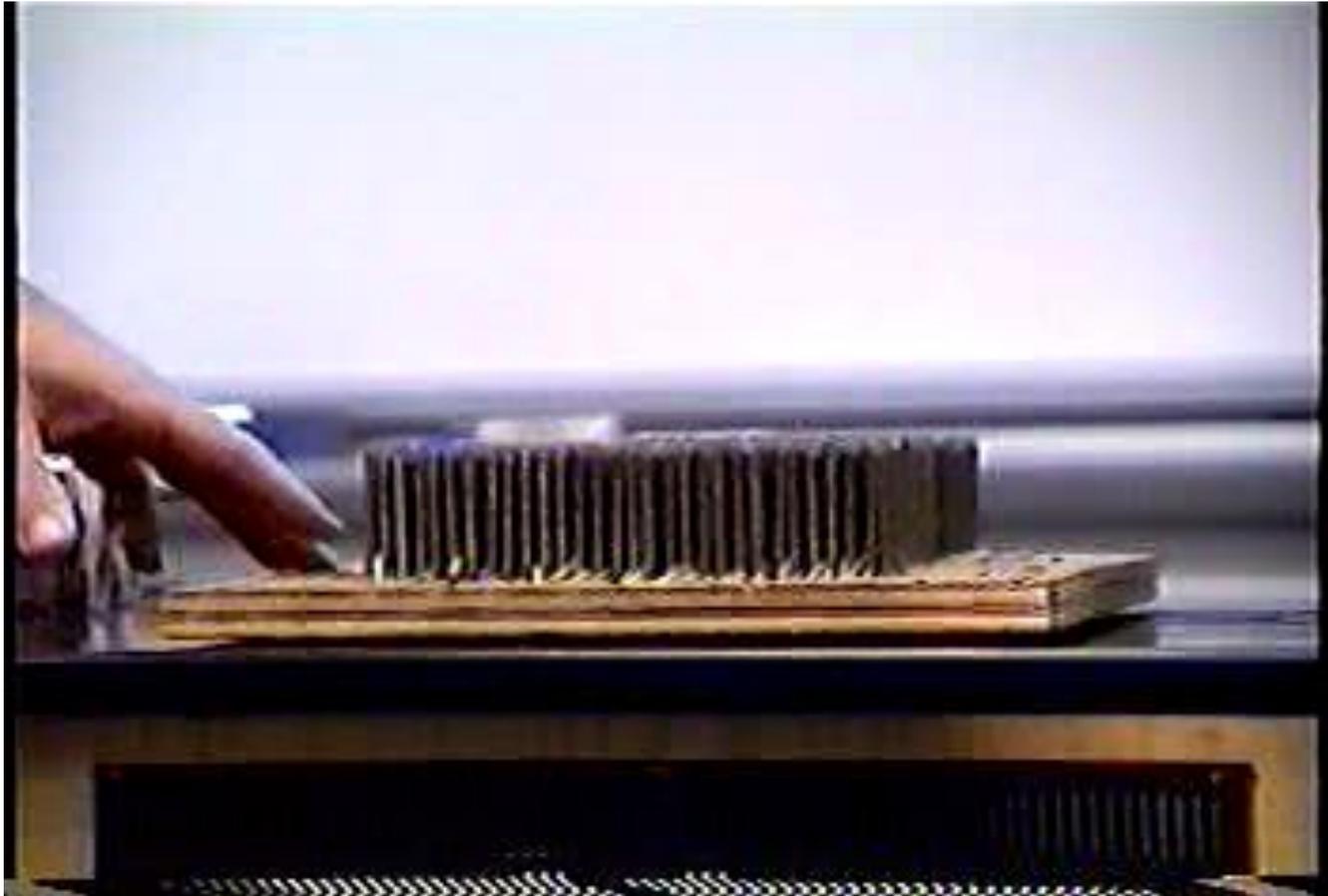


Flat & hydrophobic

Patterned & hydrophobic

References Shirlcliffe, N.J. *et al.*, *Langmuir* **21** (2005) 937-943; *Adv. Maters.* **16** (2004) 1929-1932;
27 December 2013 *J. Micromech. Microeng.* **14** (2004) 1384-1389.

Fakir's Carpet

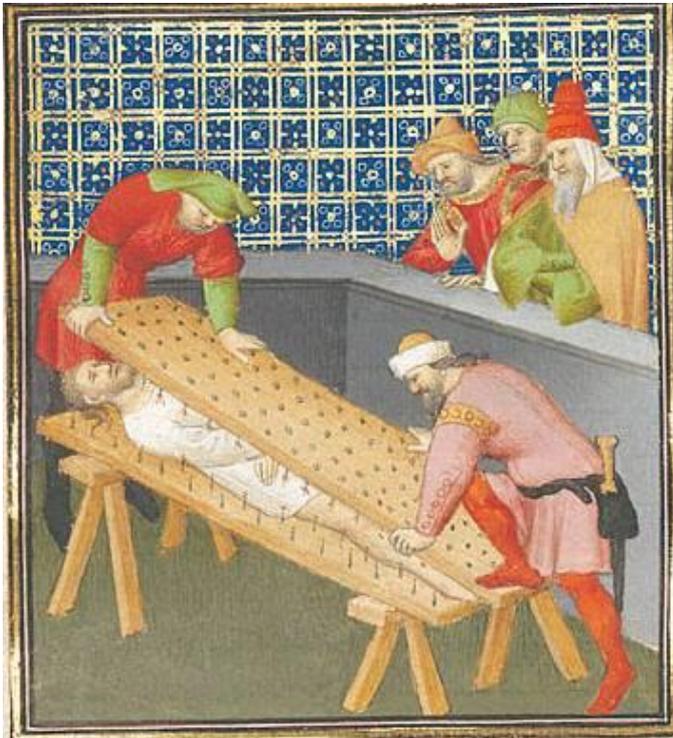


Acknowledgement: Wake Forest University

But liquid skin interacts with solid surfaces and "nails" do not need to be equally separated. A useful analogy, but it is not an exact view.

Bed of Nails

Roman consul Marcus Atilius Regulus is tortured to death by Carthaginians in about 255 BC. The illustration was painted in about 1415 in Paris.

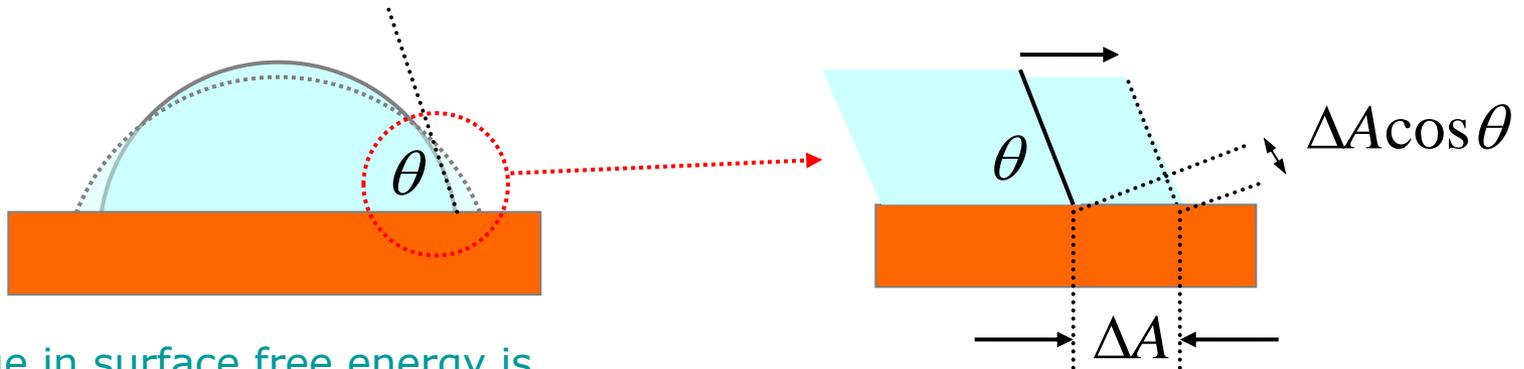


Acknowledgement: Physics, UCLA

Minimum Surface Free Energy

Young's Law – The Chemistry

What contact angle does a droplet adopt on a flat surface?



Change in surface free energy is

| | | | | |
|---|---|--|---|--|
| <p>solid-liquid gain of energy per \times substrate unit area area</p> | - | <p>solid-vapor loss of energy per \times substrate unit area area</p> | + | <p>liquid-vapor gain of energy per \times liquid-vapor unit area area</p> |
|---|---|--|---|--|

$$\Delta F(x) = (\gamma_{SL} - \gamma_{SV}) \Delta A(x) + \gamma_{LV} \Delta A(x) \cos \theta$$

Equilibrium is when $\Delta F(x) = 0 \Rightarrow$

$$\cos \theta_e = (\gamma_{SV} - \gamma_{SL}) / \gamma_{LV}$$

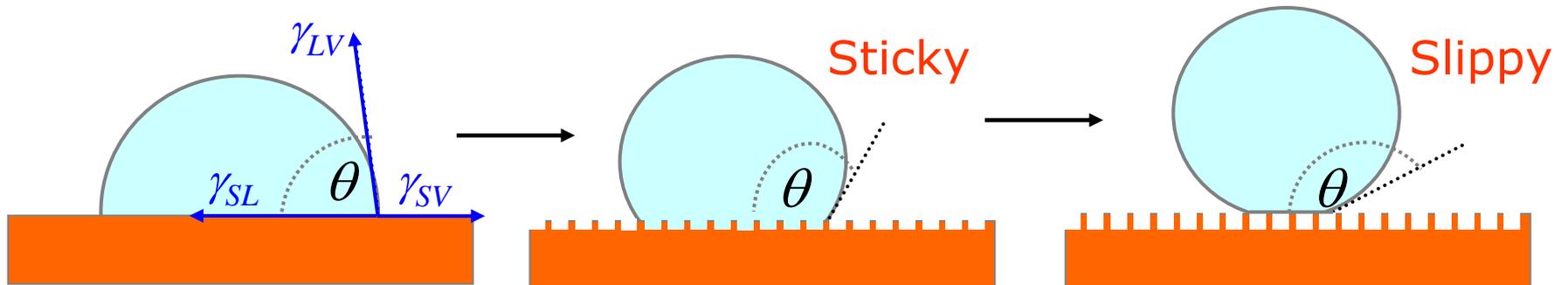
Young's Law

Same result as from resolving forces at contact line

Topography Modifies Energy Argument

Droplets that Impale and those that Skate

What contact angle does a droplet adopt on a "rough" surface?



Young's Law

Wenzel Eq.

Cassie-Baxter Eq

$$\cos \theta_e = (\gamma_{SV} - \gamma_{SL}) / \gamma_{LV}$$

$$\cos \theta_W(x) = r(x) \cos \theta_e$$

$$\cos \theta_{CB}(x) = f_s(x) \cos \theta_e - (1 - f_s(x))$$

Chemistry

Roughness

Chemistry

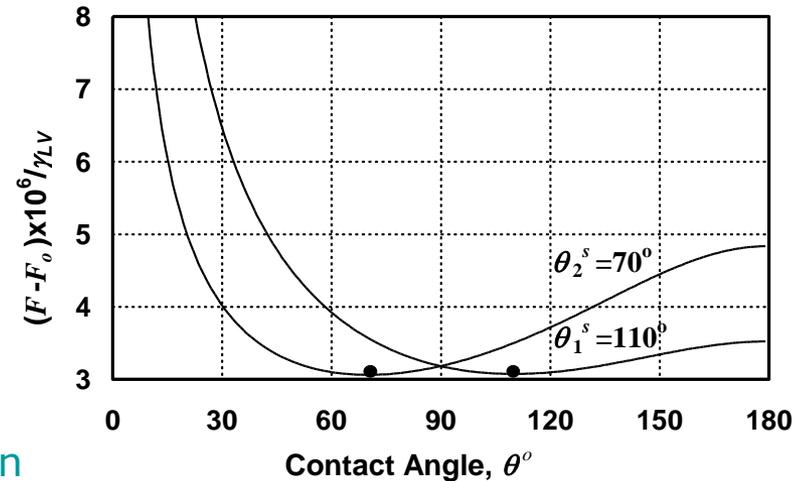
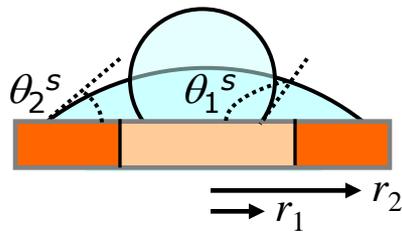
Topography

$r(x)$ = true area/planar projection at edge Young's Law θ_e $f_s(x)$ = solid surface fraction at edge

Importance of the Three-Phase Contact Line

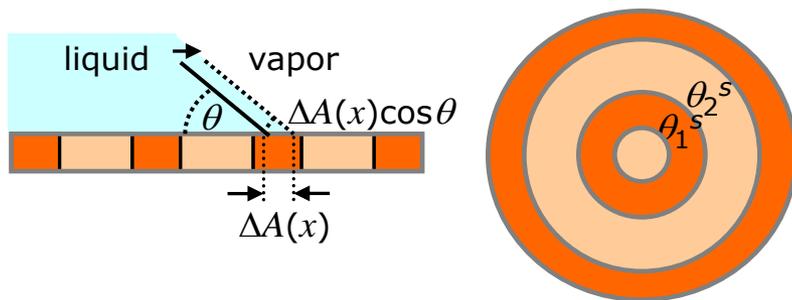
Isolated Defect Surface

Surface has $\theta_1^s = 110^\circ$, $\theta_2^s = 70^\circ$

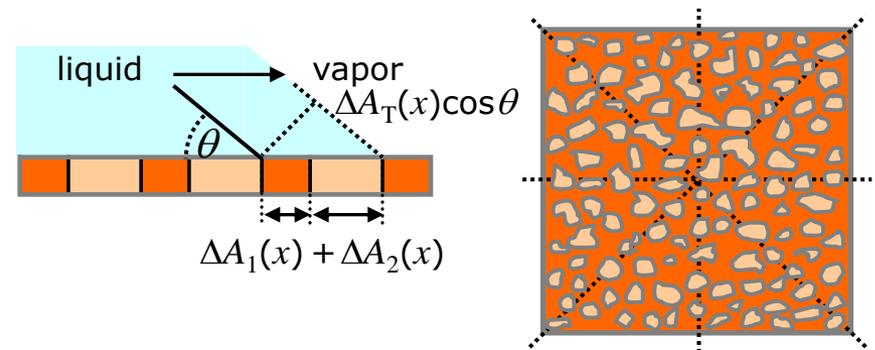


Two droplet configurations exist with min in their local surface free energy corresponding to the same droplet volume

Radial Symmetry



Random Surface



Local (not Global) Differential Parameters

Cassie-Baxter

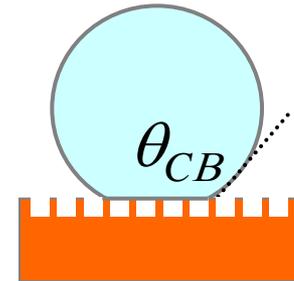
Define surface fractions: $f_i(x) = \Delta A_i(x) / (\Delta A_1(x) + \Delta A_2(x))$

$$\cos \theta_c(x) = f_1(x) \cos \theta_1 + f_2(x) \cos \theta_2$$

for a simple post-type superhydrophobic surface \Rightarrow

$$\cos \theta_{CB}(x) = f_s(x) \cos \theta_e - (1 - f_s(x))$$

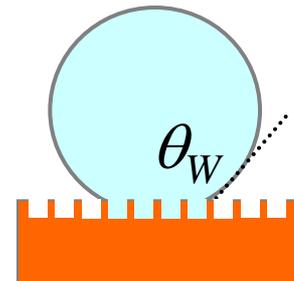
where $f_s(x)$ is the solid surface fraction and the x indicates values at the three-phase contact line ($\theta_e = \theta_e(x)$ is also local to the three-phase contact line)



Wenzel

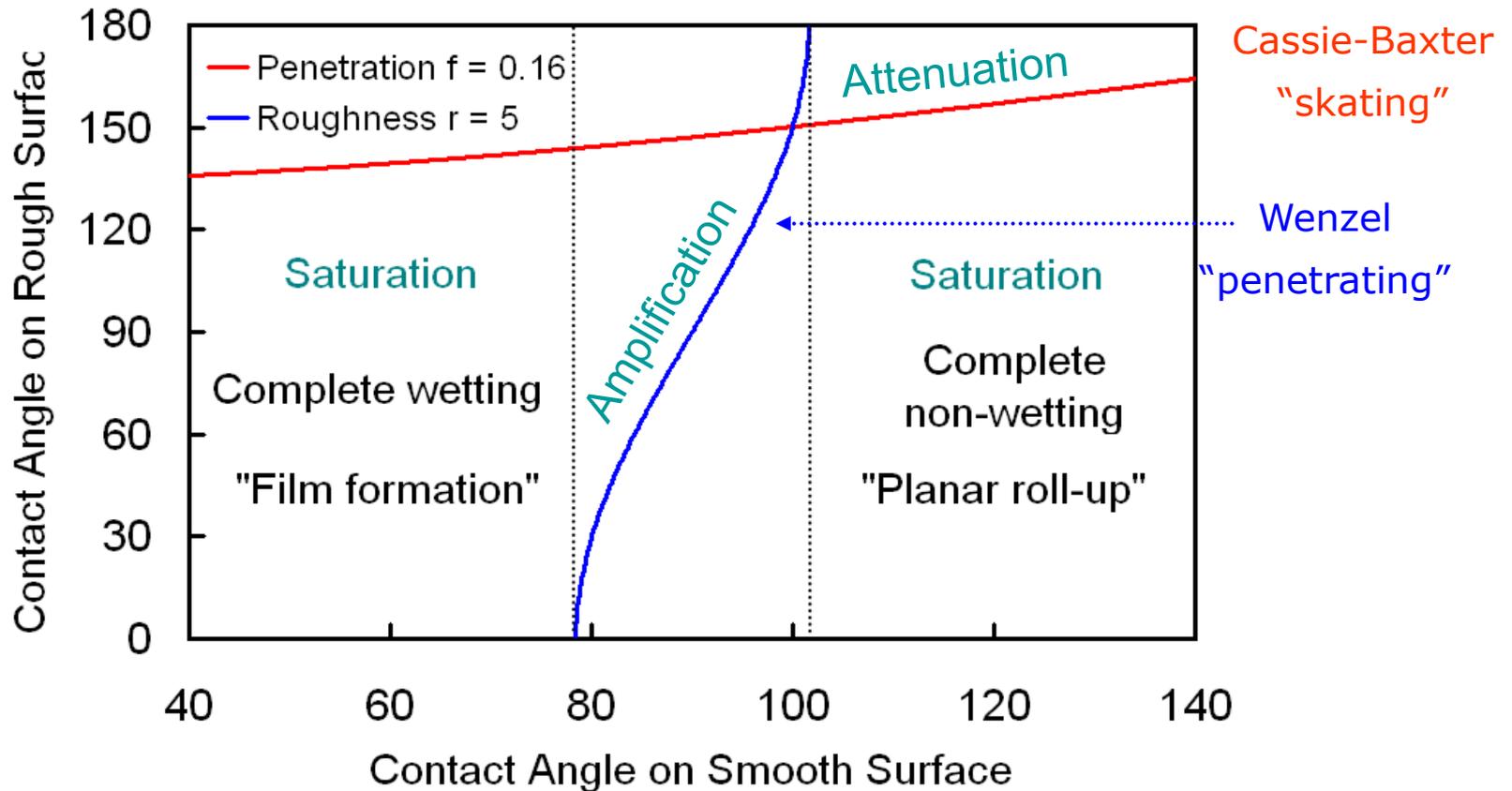
Define roughness: $r(x) = \Delta A_{wetted}(x) / \Delta A_{projected}(x)$

$$\cos \theta_W = r(x) \cos \theta_e$$

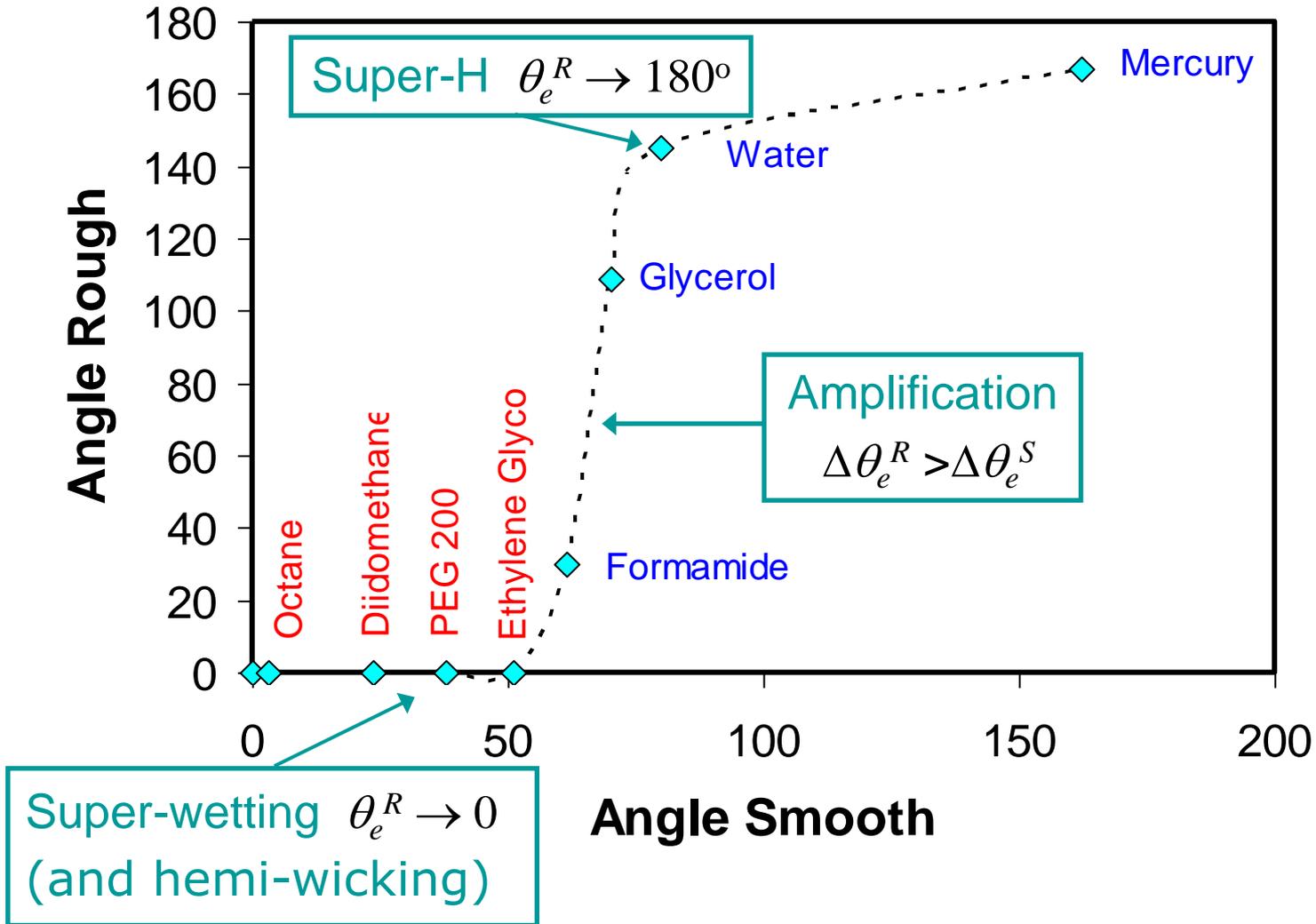


3. Amplification, Switching and Superspreading

Amplification, Attenuation, Saturation

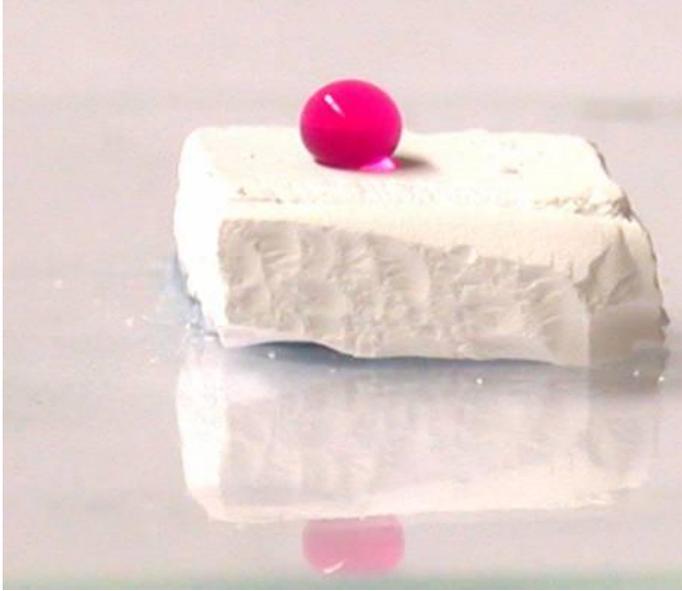


Liquids on a Superhydrophobic Surface



References McHale, G. *et al.*, *Analyst* **129** (2004) 284-287; *Langmuir* **20** (2004) 10146-10149.

Sol-Gel: Switching off Superhydrophobicity



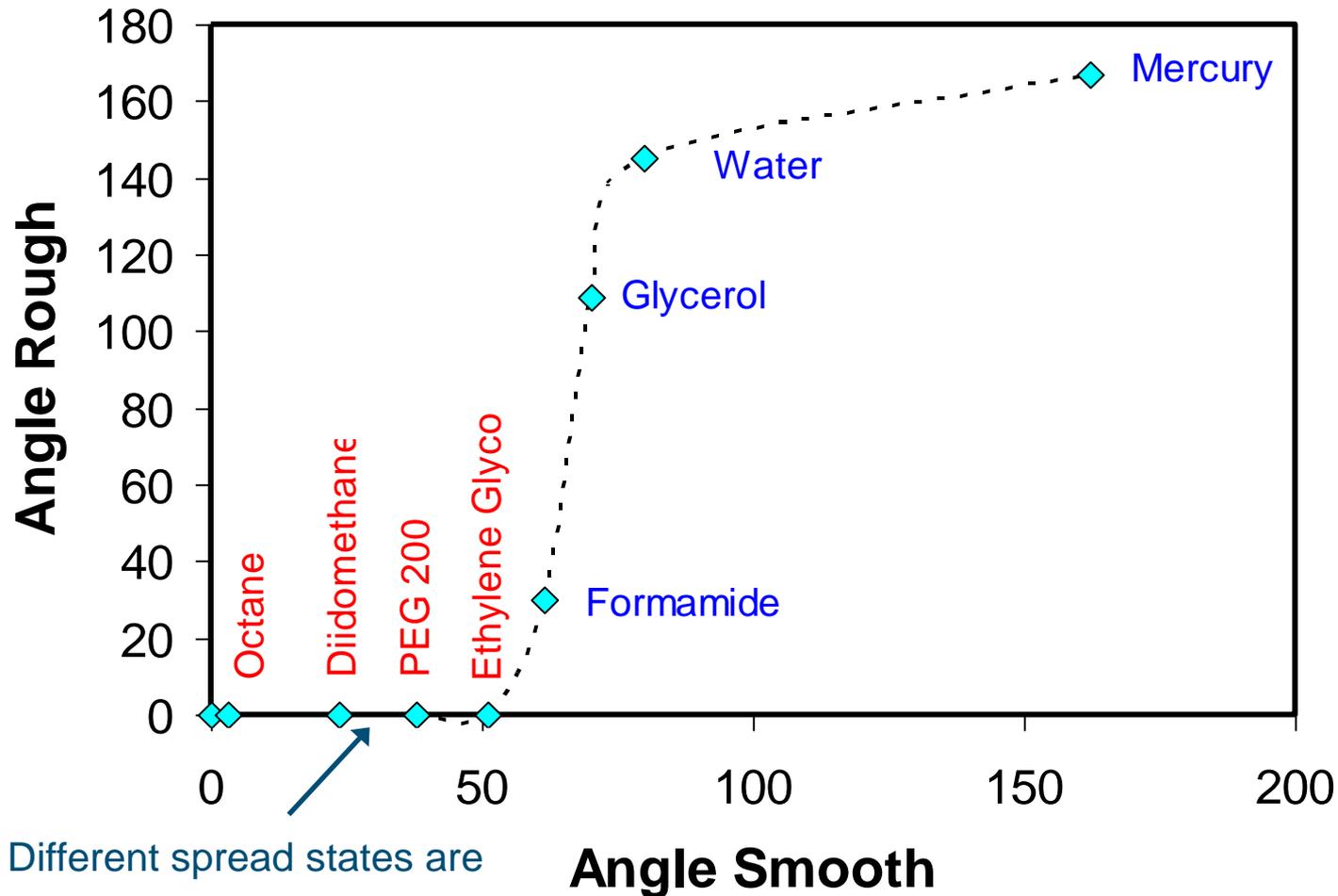
→
Foam heated
(and cooled)
prior to droplet
deposition

Mechanisms for Switching

- Temperature history of substrate
- Surface tension changes in liquid (alcohol content, surfactant, ...)
- Electrowetting

Switch could trigger a large change ⇒ Sensor based on hydrophobicity

Super-spreading

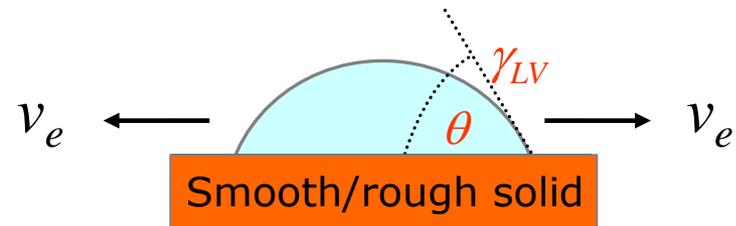


Different spread states are approached at different rates

Driving Forces for Spreading

Drop spreads radially until contact angle reaches equilibrium

Horizontally projected force $\gamma_{LV} \cos \theta$



Smooth Surface

Driving force $\sim \gamma_{LV}(\cos \theta_e^s - \cos \theta)$

Cubic drop edge speed

$$\Rightarrow v_E \propto \theta \gamma_{LV}(\theta^2 - \theta_e^{s2})$$

Wenzel Rough Surface

Driving force $\sim \gamma_{LV}(r \cos \theta_e^s - \cos \theta)$

Linear droplet edge speed

$$\Rightarrow v_E \propto \theta \gamma_{LV}((r-1) + ((\theta^2 - r\theta_e^{s2})/2))$$

Prediction : Weak roughness (or surface texture) modifies edge speed:

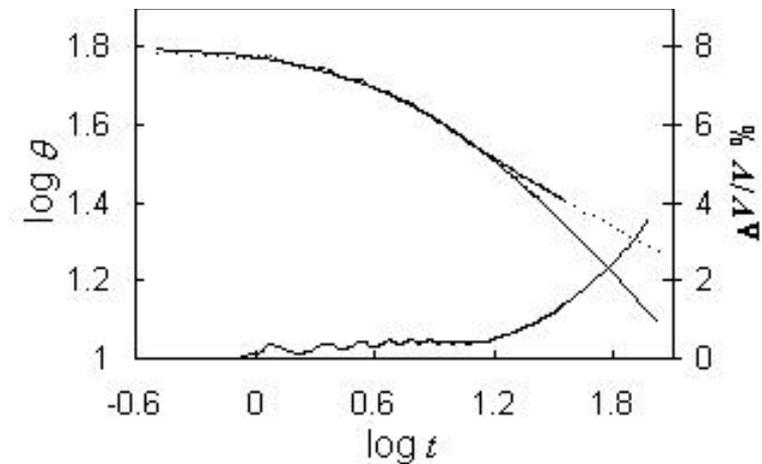
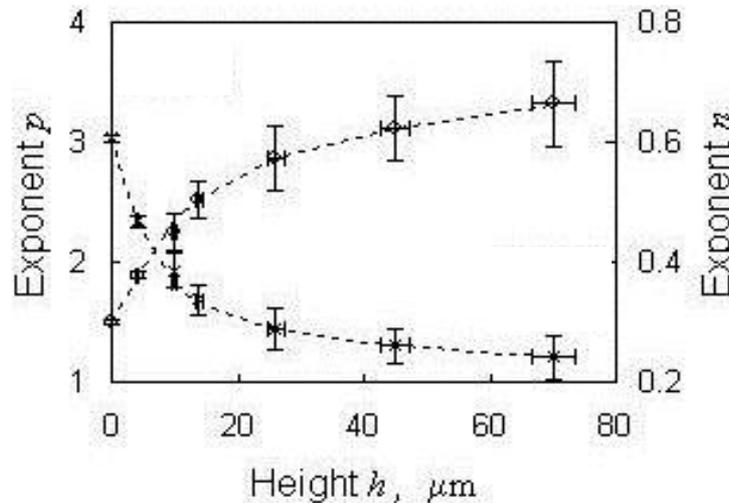
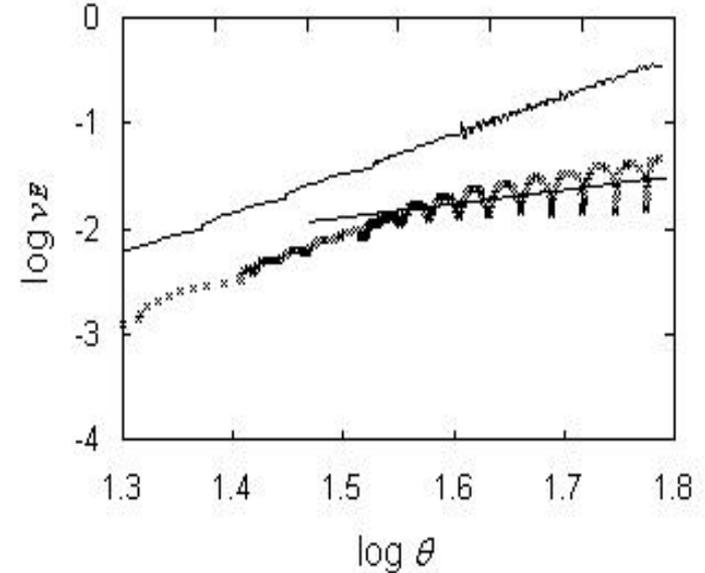
$$v_E \propto \theta(\theta^2 - \theta_e^{s2}) \quad \text{changes towards} \quad v_E \propto \theta$$

Superspreading of PDMS on Pillars

Tanner's Law exponents p and n

(cubic to linear transition)

$$v_E \propto v^* \theta^p \quad \theta \propto \left(\frac{V^{1/3}}{v^*} \right)^n \frac{1}{(t + t_o)^n}$$



References

27 December 2013

McHale, G. *et al.*, Phys. Rev. Lett. **93** (2004) art. 036102; Nature Materials. **6** (2007) 661-664.

4. Conformable Substrates

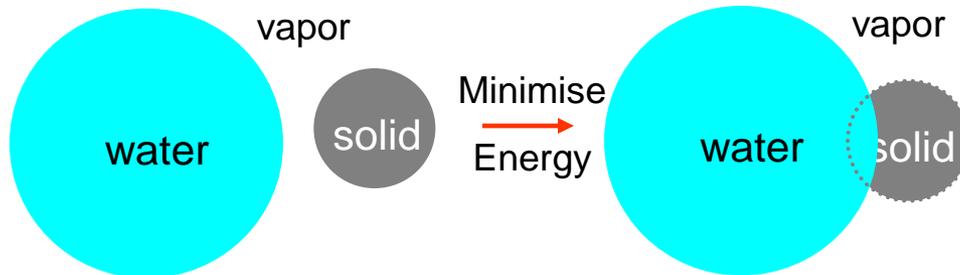
Liquid Marbles and Capillary Origami

Liquid Marbles – Assembling a Conformal Skin

Loose Surfaces

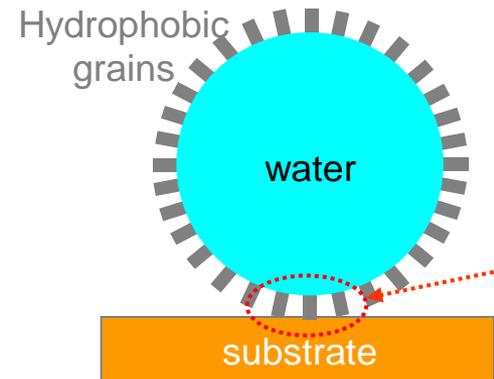
1. Grains are not fixed, but can be lifted by the liquid
2. Surface free energy favors solid grains attaching to liquid-vapor interface
3. A water droplet rolling on a hydrophobic lycopodium (or other grain/powder) becomes coated and forms a liquid marble

Hydrophobic Grains and Water



$$\Delta F = -\pi R_g^2 \gamma_{LV} (1 + \cos \theta_e) (1 + \cos \theta_e)$$

Energy is always reduced on grain attachment assuming grain is smooth ($r=1$)



Similar to pillars, but solid conformable to liquid

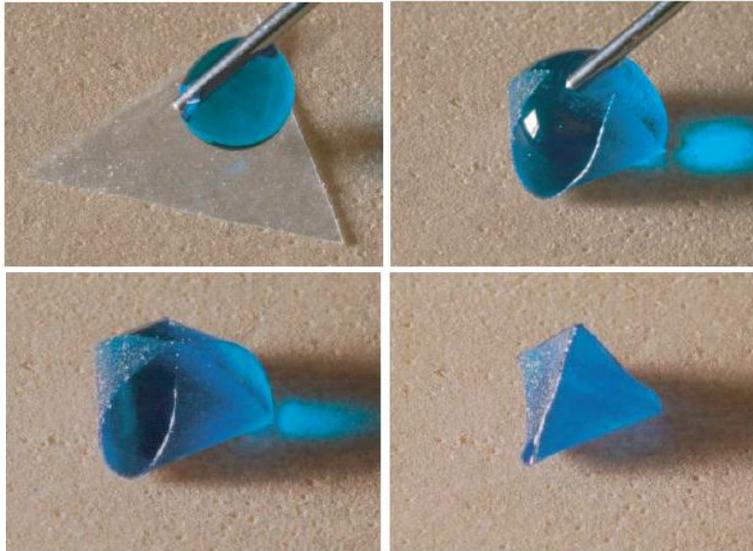


References

27 December 2013 Aussillous, P.; Quéré, D. Proc. Roy. Soc. A462 (2006) 973-999; Nature 411 (2001) 924-927. McHale, G., *et al.* 23 Langmuir (2007) 918-924. Newton, M.I., *et al.* J. Phys. D40 (2007) 20-24. McHale, G.; Newton, M.I., Soft Matter, accepted (2011).

Capillary Origami and “Hydrophilic” Teflon

Py et al's “Capillary Origami”



Water droplet contacting triangular sheet of PDMS

Acknowledgement: Py et al. Eur. Phys. J.

McCarthy's Experiment

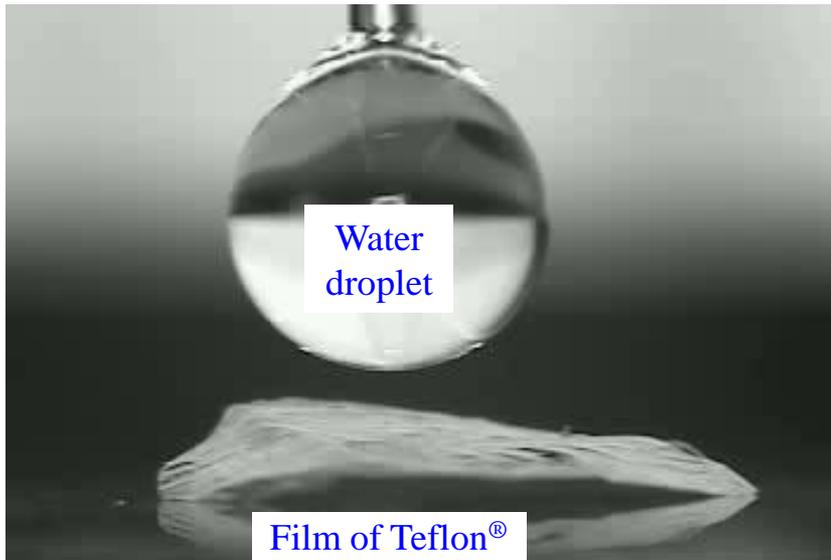


Water droplet contacting a $3.7 \mu\text{m}$ film of Teflon[®] AF2400

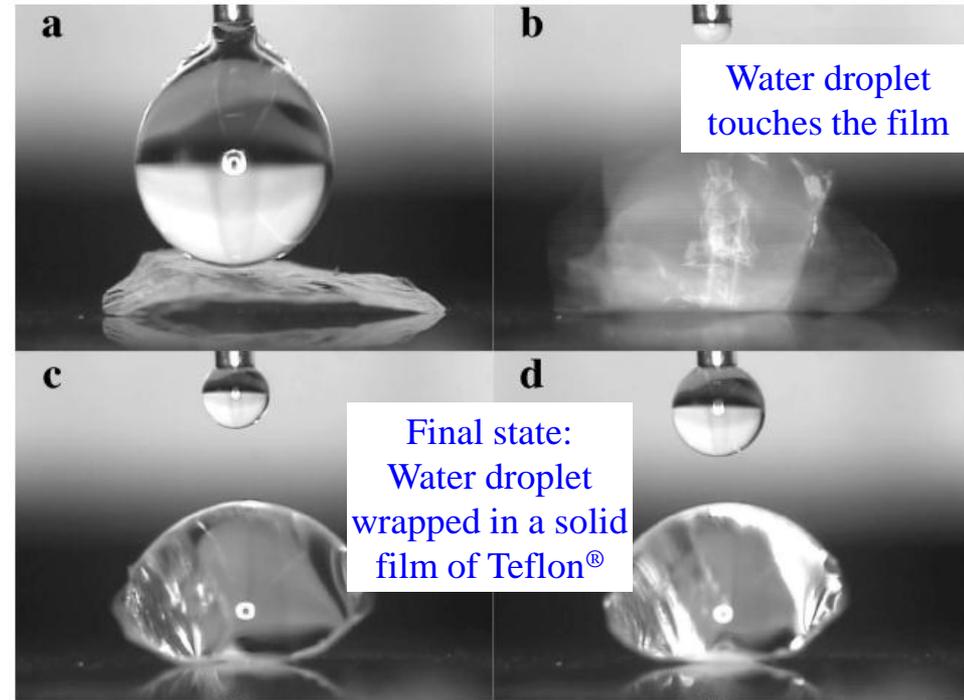
Courtesy: Prof. Tom McCarthy (UMass Amherst)

1. We all know Teflon[®] is a hydrophobic solid and gives a non-stick surface
2. Consider a thin film of Teflon contacted by a droplet of water
3. What happens?

Droplet Wrapping Video



Stills from Video



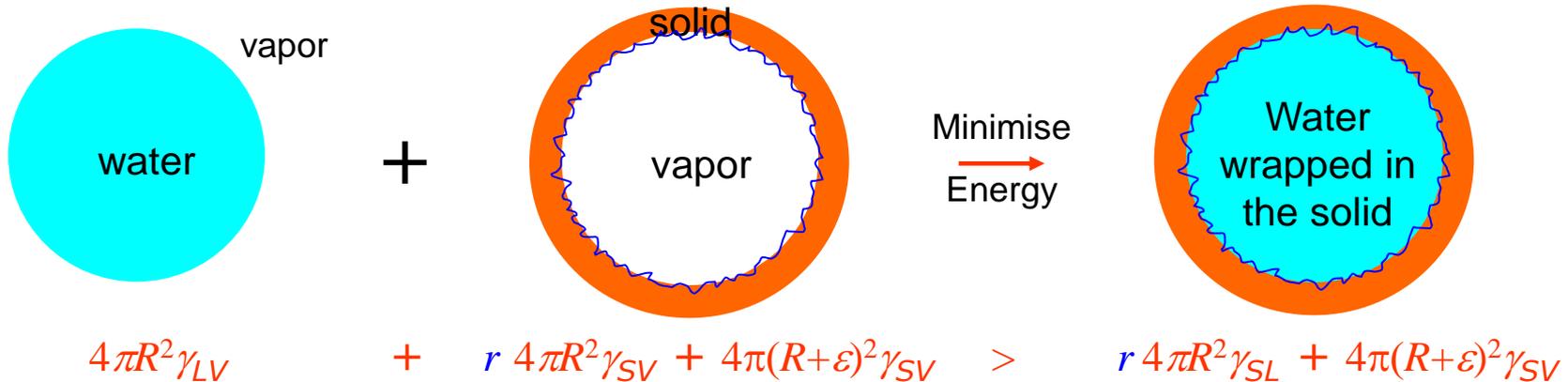
Courtesy: Prof. Tom McCarthy (UMass, Amherst)

If a droplet wraps itself up in Teflon® ... is this consistent with Teflon® being hydrophobic?

Aren't all Solids with $\theta_e < 180^\circ$ Hydrophilic?

1. Assume energy in deforming/bending solid is zero
2. Assume solid is smooth and droplet is small
3. Under these conditions surface free energy always favors solid wrapping up a droplet providing the Young's law contact angle is greater than zero

Hydrophobic Solid Shell (of thickness ε) and Water



gives $\Delta F/4\pi R^2 = r\gamma_{SL} - \gamma_{LV} - r\gamma_{SV}$ Use Young's Law $\Rightarrow = -(1 + r\cos\theta_e) < 0 \Rightarrow \theta_e < 90^\circ$ $r \rightarrow \infty$

All smooth ($r=1$) solids with Young's law $\theta_e < 180^\circ$, incl. Teflon, are absolutely hydrophilic, although those with $\theta_e > 90^\circ$ have a tendency to hydrophobicity (in a Wenzel sense)

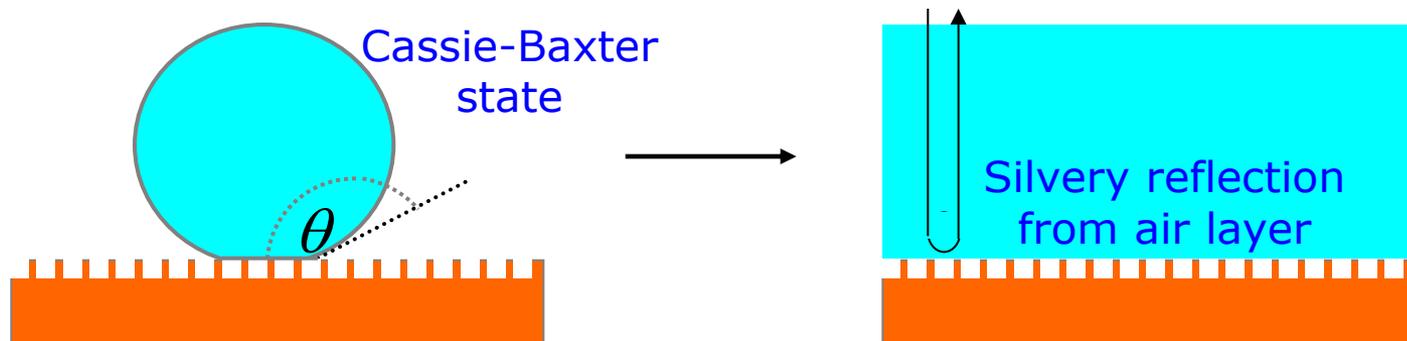
5. Immersed Superhydrophobic Surfaces

Superhydrophobicity and Plastrons

Immersed Superhydrophobic Surfaces

Provided design of features correct, penetration of water can be resisted

A silvery sheen can be seen when immersed – due to surface retained layer of air.



Plastron Respiration

Insect physiologists have studied immersed S/H surfaces since 1940's

Silvery sheen from air layer indicates an air-water interface, e.g diving spider - acts as a gas exchange membrane to extract oxygen from water and remove carbon dioxide from insect

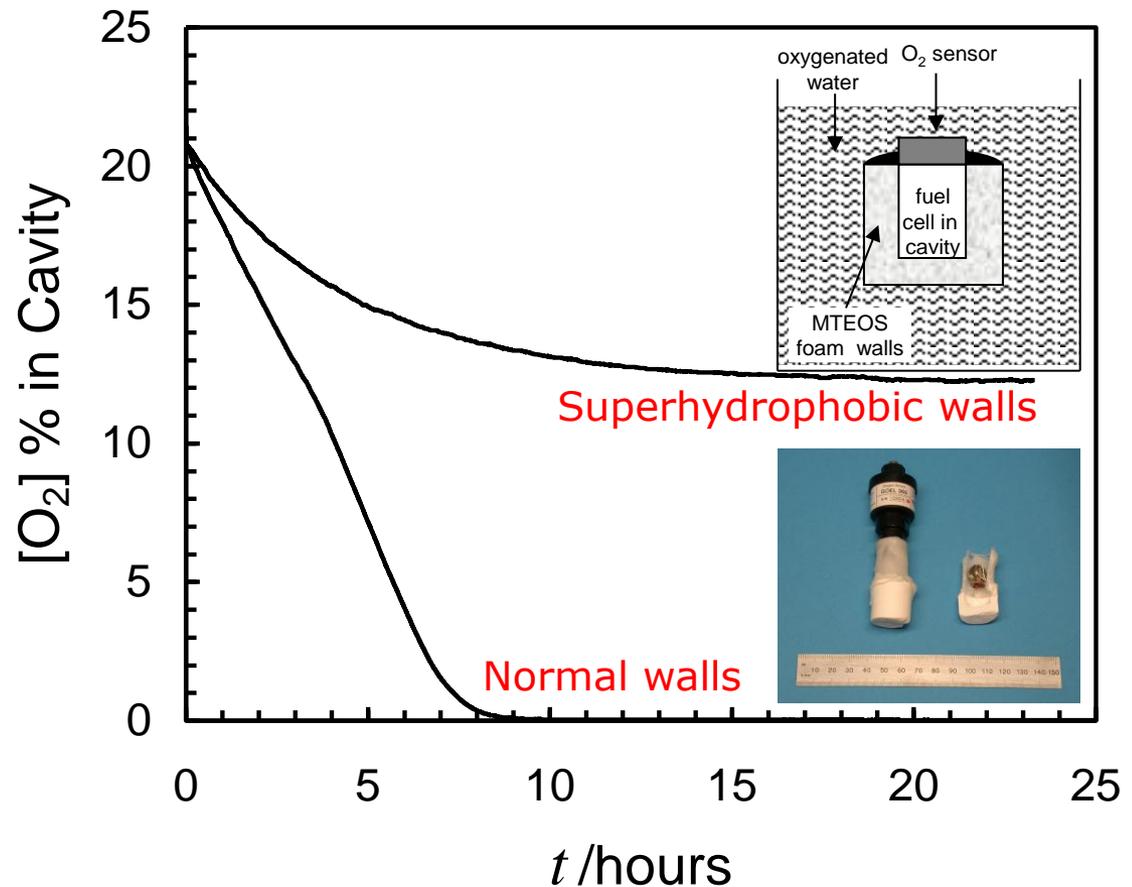
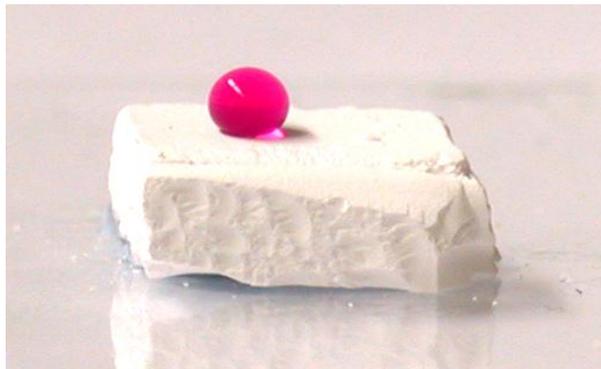
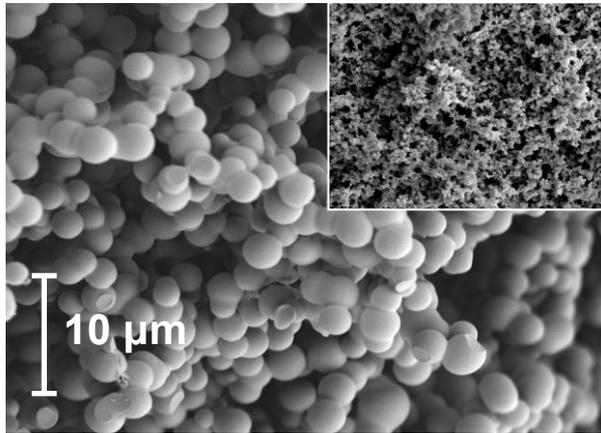


Microcosmos ©Allied Films Ltd (1996)

Plastron/Underwater Respiration Biomimic

Intrinsically superhydrophobic MTEOS organo-silica sol-gel foam surfaces

Structure controllable from nano- to macro-porous



References Shirtcliffe *et al.*, *Langmuir* **19** (2003) 5626-5631; *Appl. Phys. Lett.* **89** (2006) art. 104106.

27 December 2013 Flynn, M.R.; Bush, J.W.M., *J. Fluid Mech.* **608** (2008) 275-296.

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Superhydrophobicity: Plastron Respiration

Similar to super gas exchange membranes

Edward Cussler

Underwater Breathing:

BBC Radio 4 Broadcast

Edward Cussler, Professor of Chemical
Engineering (University of Minnesota)

Speaking 9th February 2006

Flow in Pipes with Superhydrophobic Walls

Concept



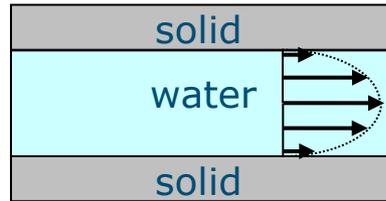
Experiment

Forced flow through small-bore Cu tubes

Electron microscope images of hydrophobic nano-ribbon ($1\mu\text{m} \times 100\text{nm} \times 6\text{nm}$) decorated internal copper surfaces of tubes (0.876 mm radii).

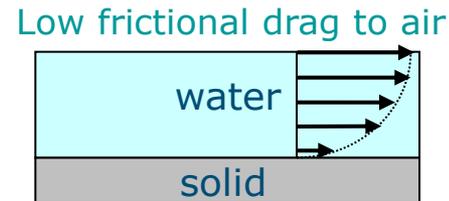
Side-profile optical images of droplets of b) water, and c) glycerol on surface shown in a) the original surface is shown in d)

Closed-channel



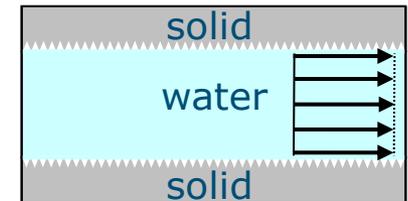
Two walls cause frictional drag

Open-channel

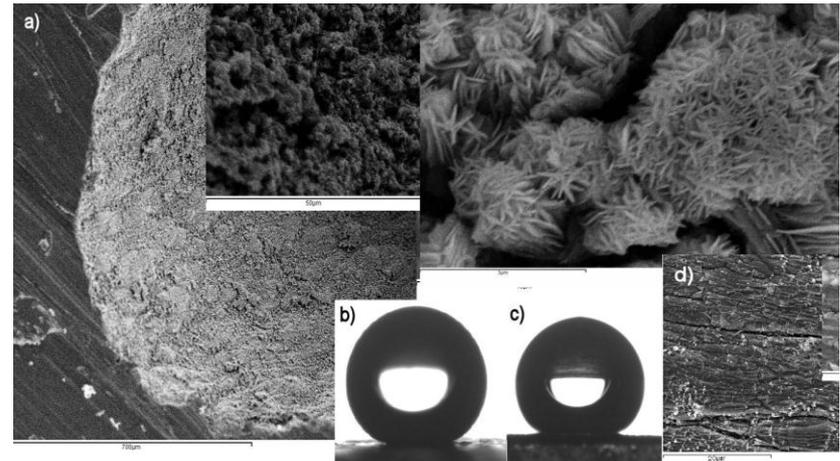


High frictional drag to solid

Super-channel



Walls appear as cushions of air



Reference

Shirtcliffe, N.J., et al.. ACS Appl. Mater. Interf., 1 (2009) 1316-1323.

27 December 2013

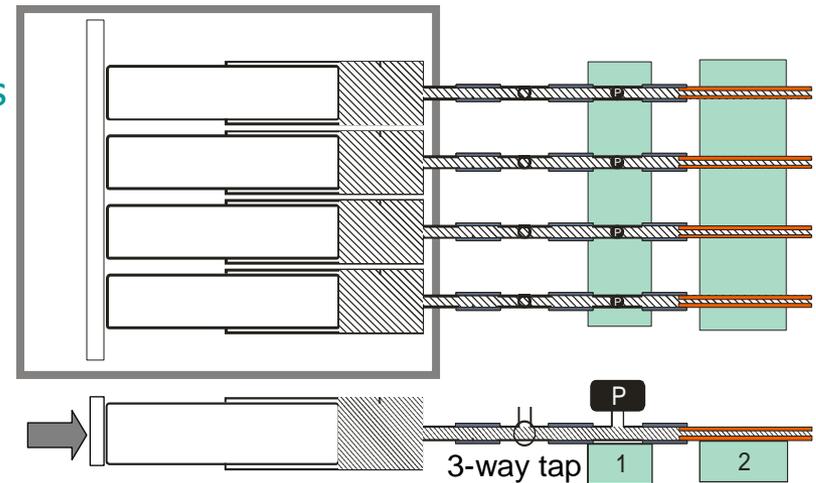
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Flow Experiments

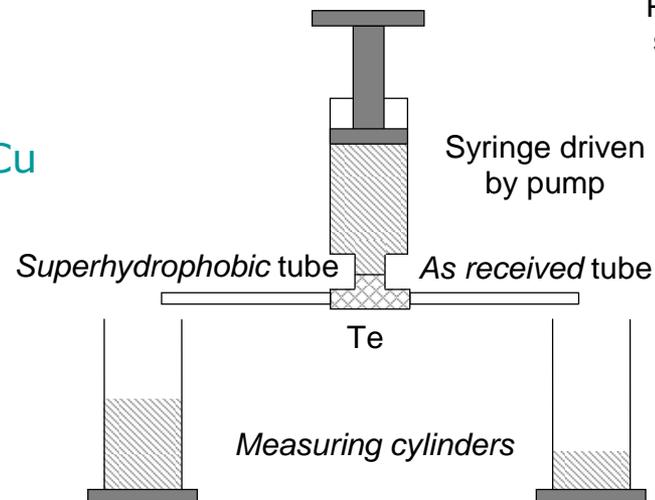
Quantitative Experiment

1. 4 parallel tubes with 4 surface finishes
2. Cu, hydrophobic Cu, nanoribbon Cu and hydrophobic nanoribbon Cu
3. Syringe pump to force flow in all 4
4. Measure pressure drop across each



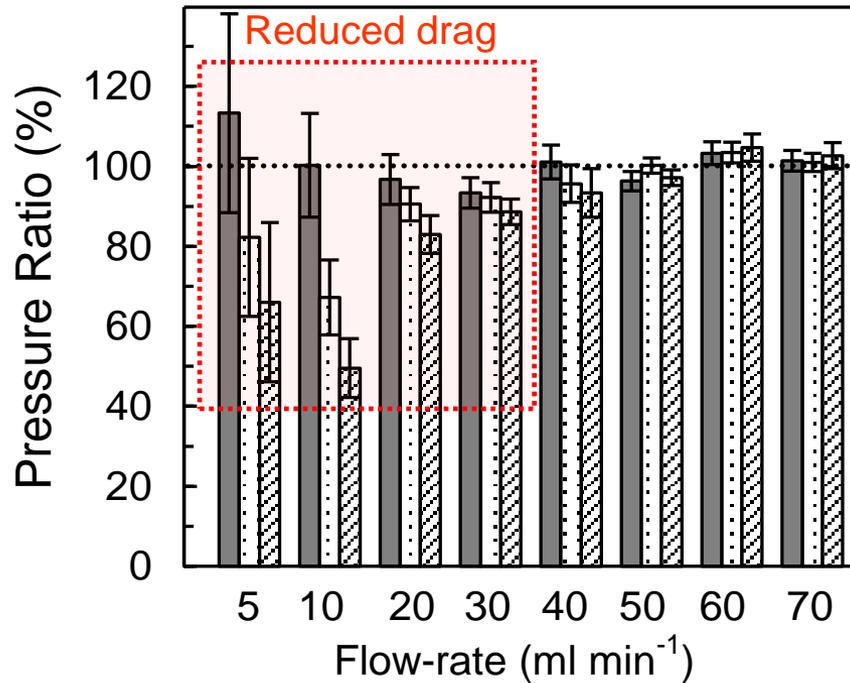
Visualization Experiment

1. 2 tubes in T-arrangement
2. Cu and hydrophobic nanoribbon Cu
3. Syringe pump to force flow
4. Outlet volumes collected

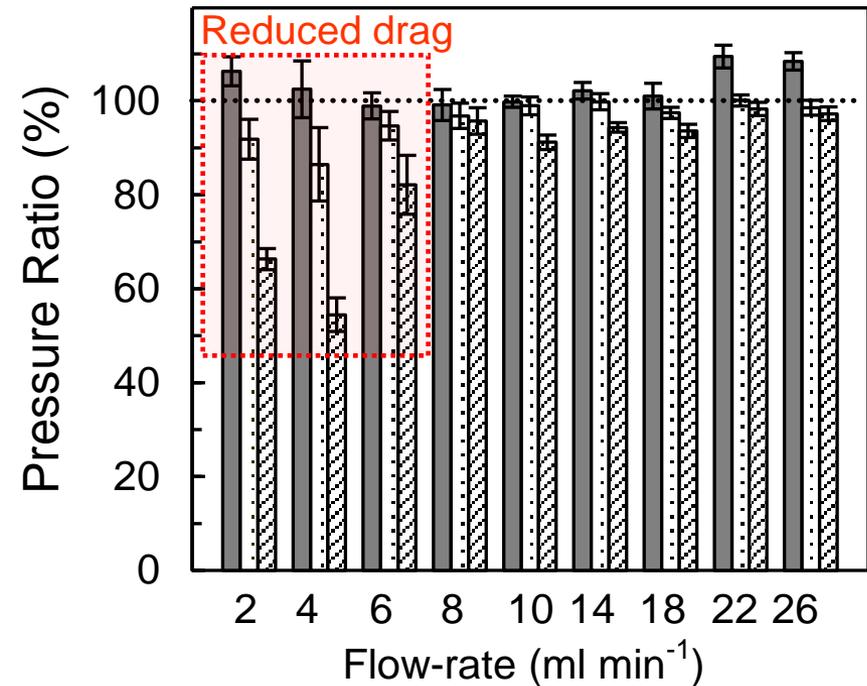


Quantitative Results

Water



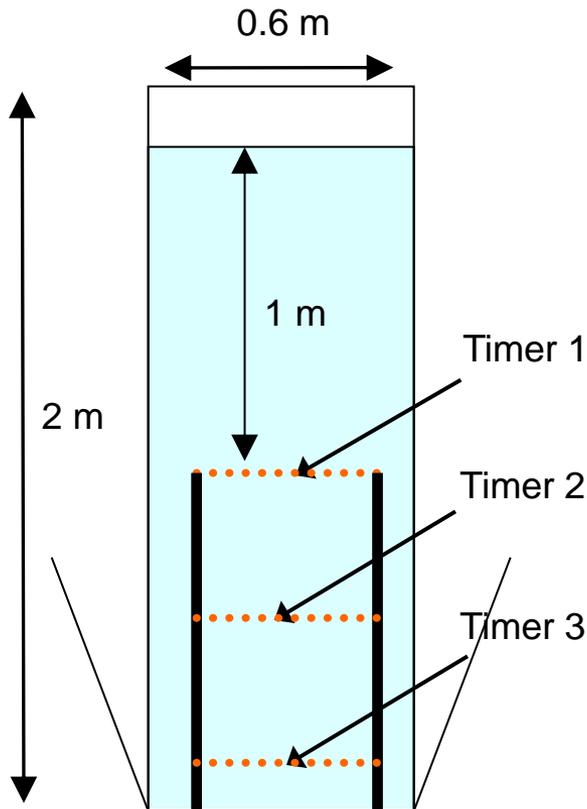
Water-Glycerol (50%)



Copper tubes with superhydrophobic inner surfaces show significantly increased flow-rates

Terminal Velocity Experiments

In the presence of a fluid, a falling object eventually reaches a terminal velocity. Textbooks tell us that in water the terminal velocity does not depend on the surface chemistry But is that true?

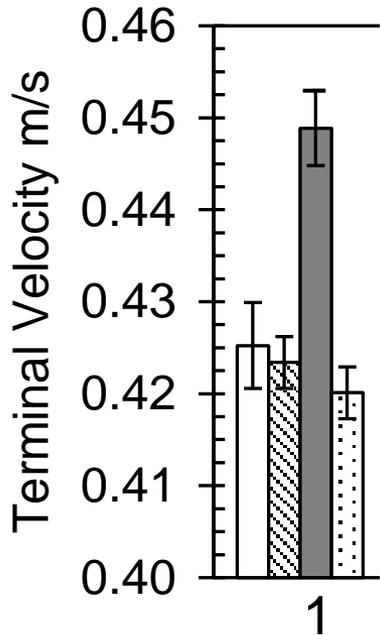


Solid sphere
Plastron bearing sphere
Same sphere



Terminal Velocity Results

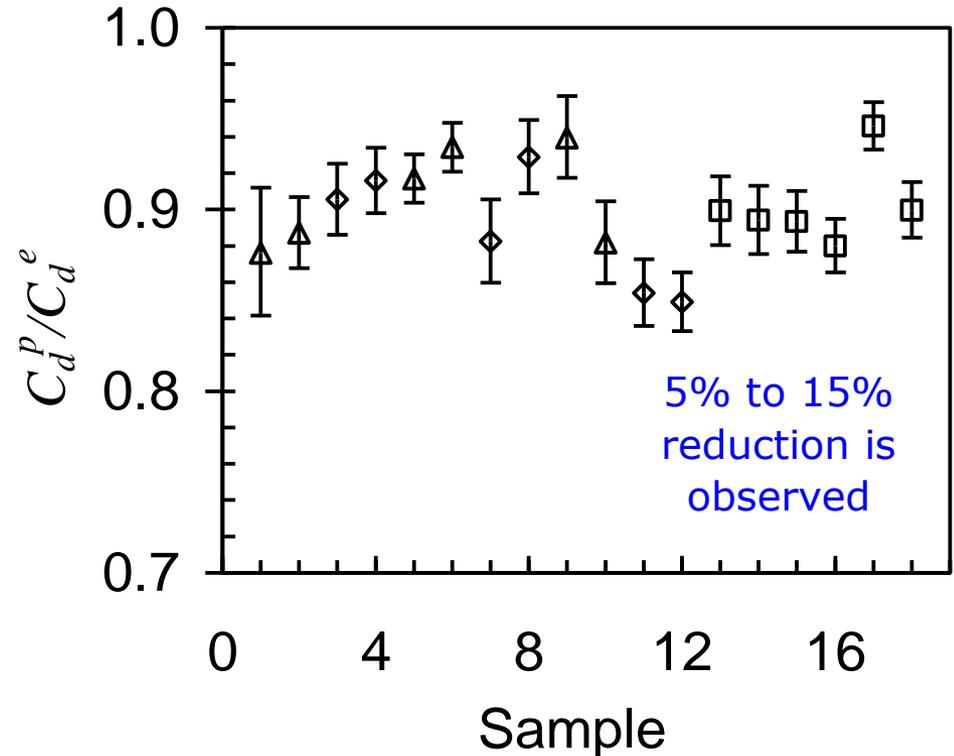
Results for 1-inch Diameter Sphere



Sequence of Four Bars

1. Blank surface
2. Sieved sand surface
3. (Super) Hydrophobic sand
4. Hydrophobic sand with ethanol pre-treatment to prevent plastron

Reduction in Drag Coefficient

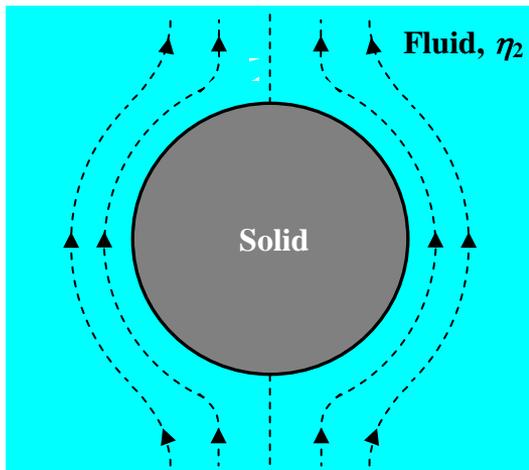


Superhydrophobicity alone is not enough. Also need a plastron to persist to achieve drag reduction

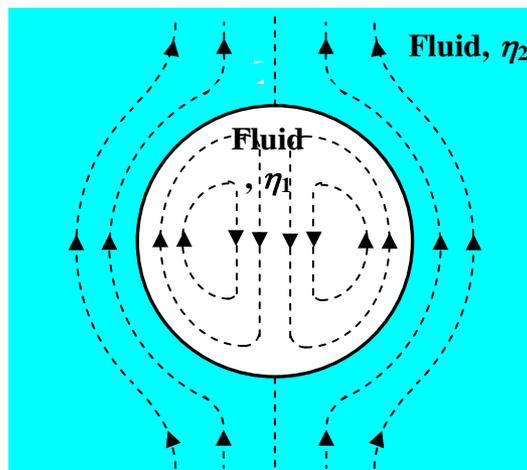
Drag Reduction – Boundary Conditions

Fundamental boundary condition is not "no-slip", but is continuity of shear stress
Well-known drag reduction effects for gas bubbles with non-rigid interfaces in water

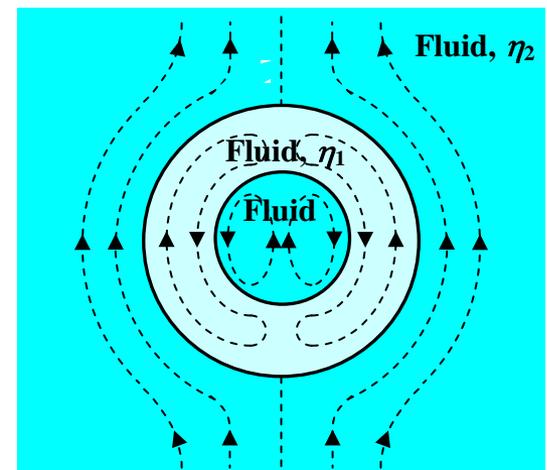
Stokes Drag (Low Re)



Hadamard-Rybczynski



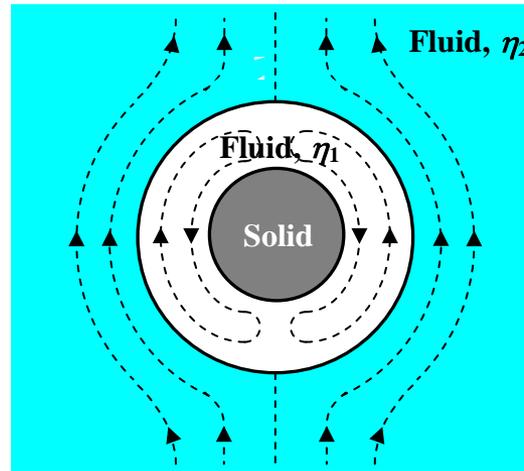
Encapsulated Droplet



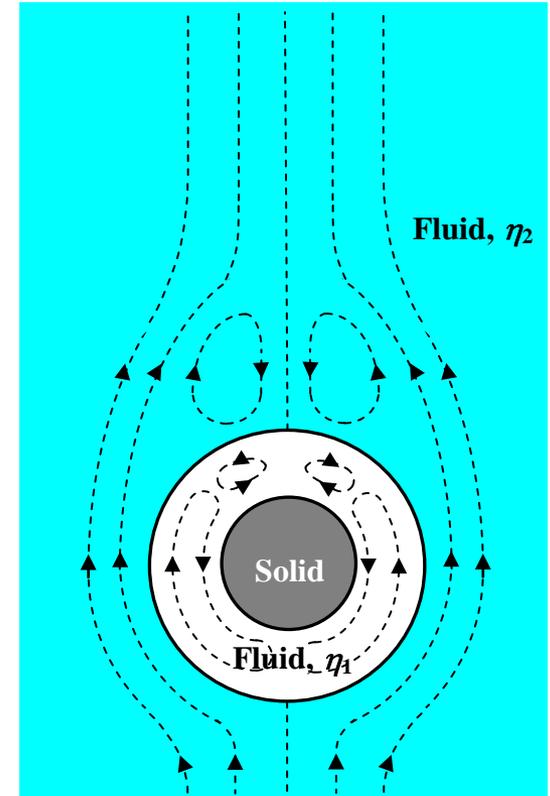
Hadamard-Rybczynski drag is 25% less than Stokes drag

Low and Intermediate Re numbers?

At low Re , a plastron/air layer may cause a Hadamard-Rybczynski effect and reduce drag



At intermediate Re , an air layer is likely to alter flow patterns and modify wake separation – possibly reducing drag

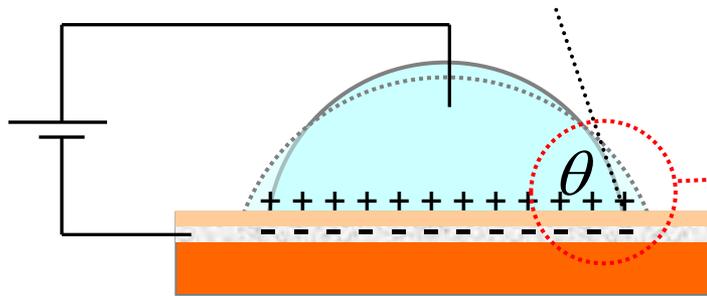
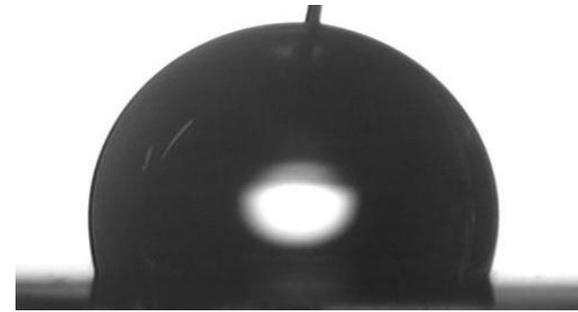
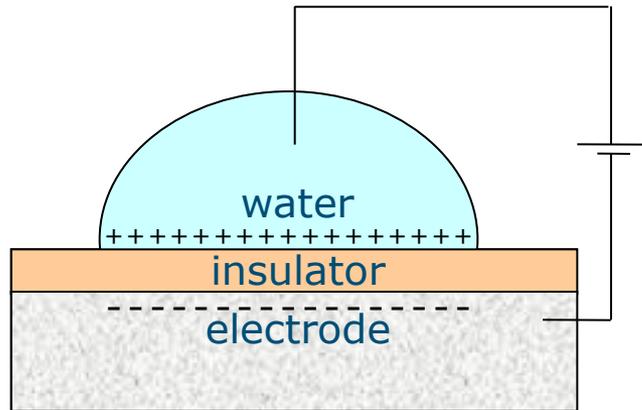


A persistent plastron/air layer is needed to achieve drag reduction

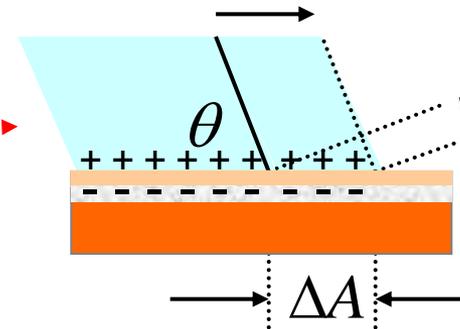
6. Dielectrowetting

Led by Prof. Carl Brown

The Principle of Electrowetting-on-Dielectric



Insulator
Metal
Substrate



$$\Delta A \cos \theta$$

$$\Delta W_e = \Delta A c V^2 / 2$$

⇒

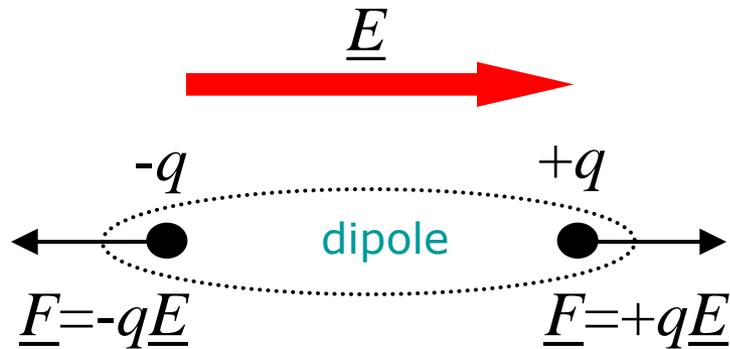
$$\cos \theta_e(V) = \cos \theta_e(0) + \frac{\epsilon_r \epsilon_0 V^2}{2d\gamma_{LV}}$$

EWOD Modified
Young's Law

Dielectrophoretic (DEP) Forces

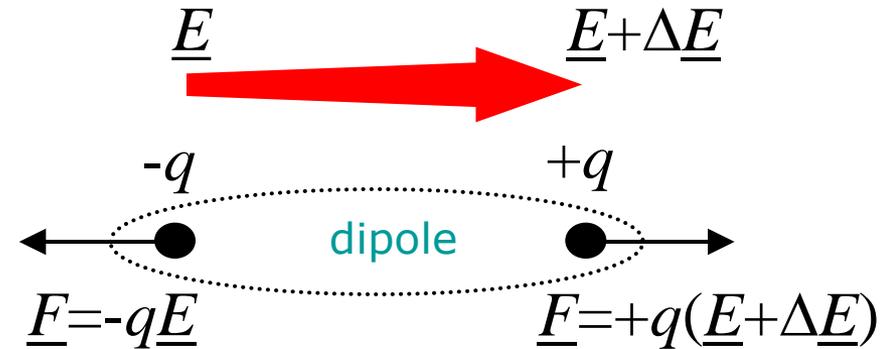
Force on a dipole in a dielectric material caused by a non-uniform electric field

Uniform Electric Field Applied



Zero net force on dipole

Non-Uniform Electric Field Applied



Net force on dipole = $+q\Delta\underline{E}$

In a dielectric liquid a non-uniform electric field causes liquid motion

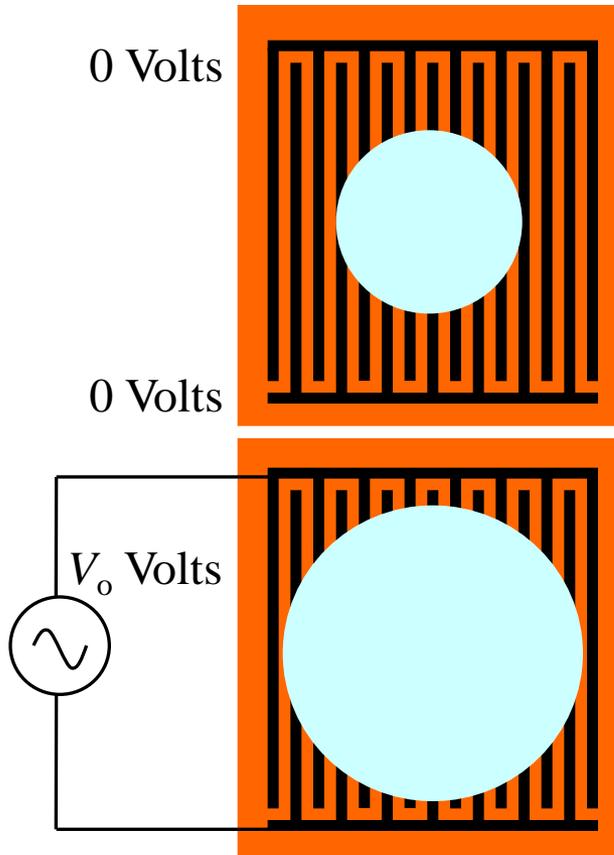
L-DEP Comparison to Electrowetting-on-Dielectric (EWOD)

1. L-DEP acts on the bulk material, but EWOD acts at the contact line
2. L-DEP uses dielectric liquids, but EWOD uses conducting liquids
3. L-DEP does not require electrical contact, but EWOD does require a contact

L-DEP Driven Spreading

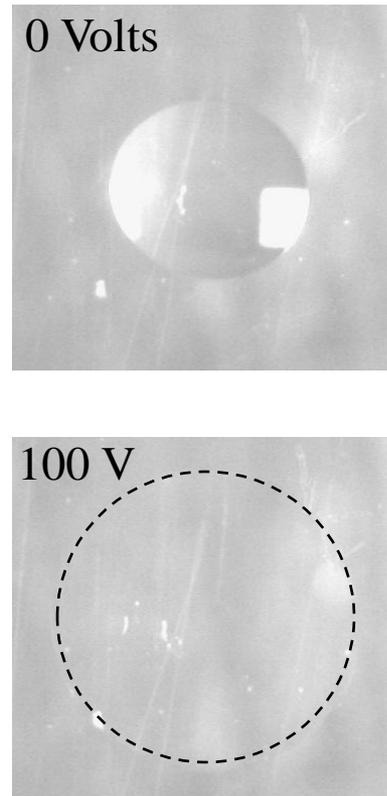
Interdigital Transducers

Top view



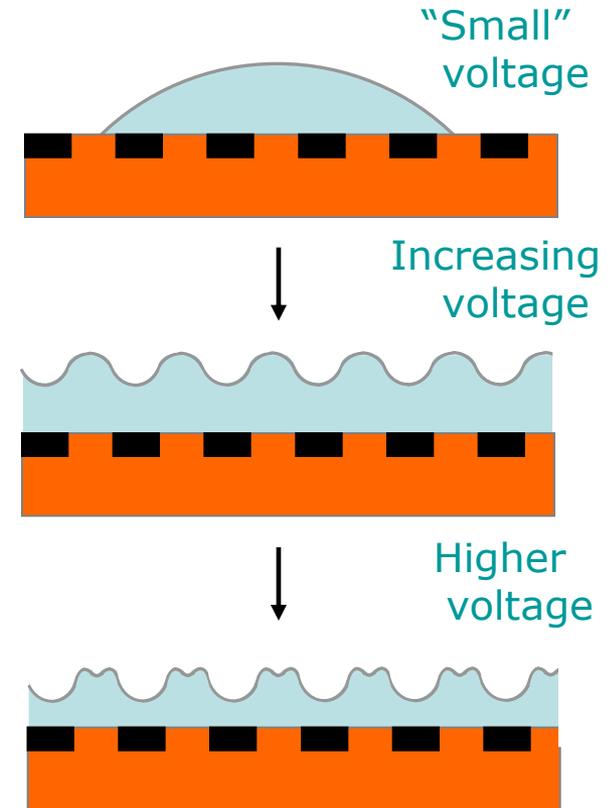
Castor Oil Droplet

Top view*

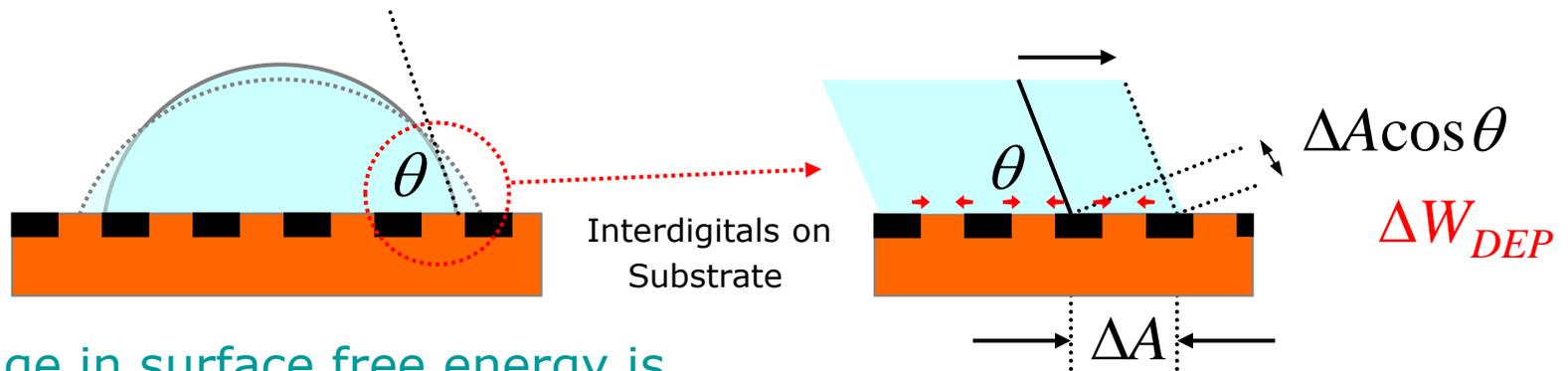


Droplet to Wrinkled Film

Side view



Surface Free Energy: Droplet Spreading



Change in surface free energy is

$$\Delta F = (\gamma_{SL} - \gamma_{SV}) \Delta A + \gamma_{LV} \Delta A \cos \theta - \Delta W_{DEP}$$

Exponential field decay into liquid, $E_z = \alpha V \exp(-\alpha z) \Rightarrow$ pen. depth $z \sim 1/2\alpha$.

L-DEP energy change is, $\Delta W_{DEP} = \epsilon_r \epsilon_0 \alpha^2 V^2 \Delta A / 4\alpha$, assuming a thick droplet

Equilibrium is when $\Delta F = 0$



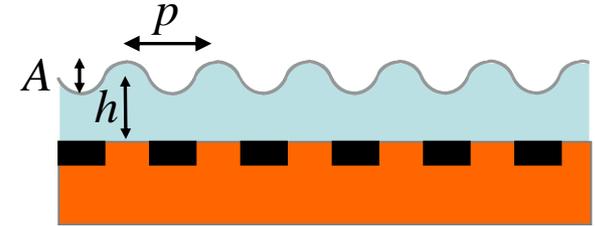
$$\cos \theta_e(V) = \cos \theta_e(0) + \epsilon_r \epsilon_0 \alpha V^2 / 4\gamma_{LV}$$

L-DEP Modified
Young's Law

Surface Free Energy: Sinusoidal Wrinkles

Thin Droplet Case (Far-field Wrinkles)

1. Electric field penetrates to upper liquid-air interface
2. Deformation of liquid-air interface can change surface energy
3. Redistribution of liquid in a pattern following "smoothed" field of IDT alters capacitive energy



Additional surface area:

$$\Delta A_{LV} = \pi^2 A^2 / 2p$$

Decrease in capacitive energy:

$$\Delta W_{DEP} = \Delta C V(z)^2 / 2 \approx \Delta C V^2 \exp(-2\pi h/p) / 2$$

Capacitance is a function of h/p and scales with $\epsilon_r \epsilon_0$, i.e. $C = \epsilon_r \epsilon_0 f(h/p)$

Change in capacitance is:

$$\Delta C = (\epsilon_r \epsilon_0 A/p) [df/du]_{u=h/p}$$

Minimizing energy with respect to changes in amplitude A ,

⇒

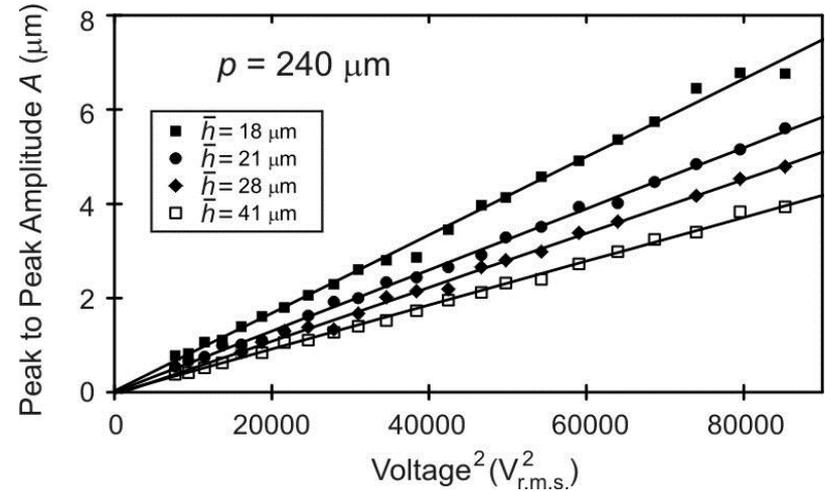
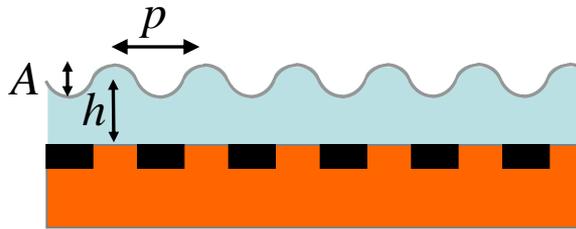
$$A \propto \epsilon_r \epsilon_0 V^2 \exp(-2\pi h/p) / 4 \gamma_{LV}$$

Amplitude Scaling
Law*

**Full solution of Maxwell's equation gives same results*

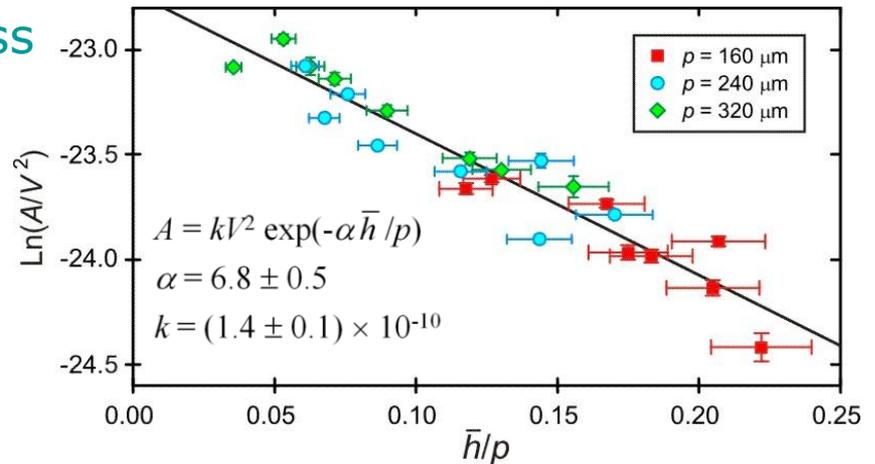
Experimental: Observed Scaling Laws

10 kHz sinewave, 1-decanol oil
 $p=160, 240$ and $320 \mu\text{m}$



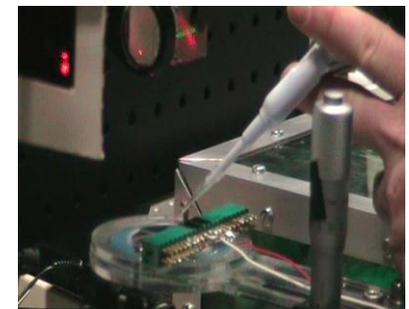
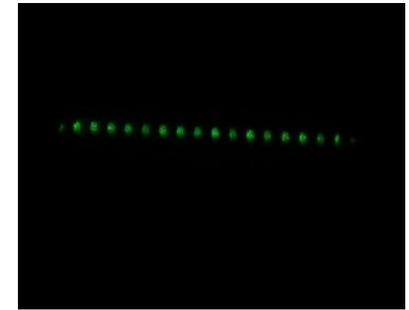
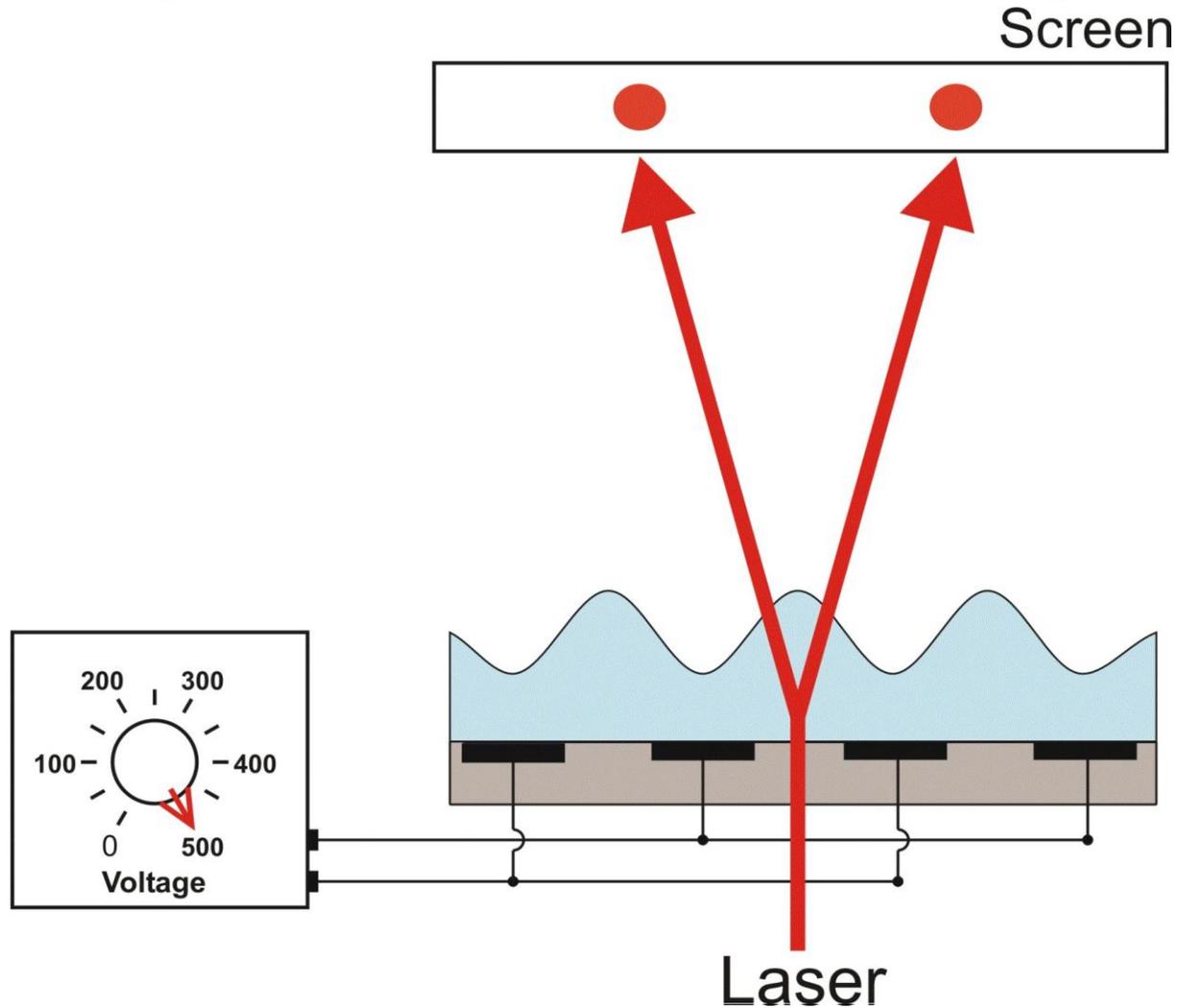
Scaling of amplitude with thickness to electrode periodicity:

$$A = kV^2 \exp(-2\pi h/p)$$



Reference Brown, C.V.; et al., Appl. Phys. Lett. 97 (2010) art. 242904.

Programmable Phase Grating



Reference Brown, C.V.; *et al.*, Nature Photonics. 3 (2009) 403-405.

27 December 2013

Summary and Conclusions

1. Topography enhancement of wetting is rich in effects

- Superhydrophobicity
- Switching
- Superwetting and superspreading

2. Substrates can be shaped by the liquid

- Hydrophobic, but adhesive
- Liquid Marble
- Capillary Origami

3. Underwater superhydrophobicity offers functional properties

- Plastron/underwater respiration without gills
- Suppression/reduction of adsorption
- Reduction of drag

4. Dielectrowetting

- Control of wetting of oils
- Shaping surfaces
- Creating optical effects

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Collaborators

| | |
|-----------|---|
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EP/C509161/1 – Extreme soil water repellence

EP/D500826/1 & EP/E043097/1 – Slip & drag reduction

EP/E063489/1 – Exploiting the solid-liquid interface

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EU COST Action D19 - Chemistry at the nanoscale

EU COST Action P21 - Physics of droplets

